

## Optical Phase Conjugation with Less Than a Photon per Degree of Freedom

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We demonstrate experimentally that optical phase conjugation can be used to focus light through strongly scattering media even when far less than a photon per optical degree of freedom is detected. We found that the best achievable intensity contrast is equal to the total number of detected photons, as long as the resolution of the system is high enough. Our results demonstrate that phase conjugation can be used even when the photon budget is extremely low, such as in high-speed focusing through dynamic media or imaging deep inside tissue.

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Biological tissue strongly scatters light, traditionally limiting the depth of optical imaging to within one millimeter. Recently, methods where phase conjugate fields “time reverse” the effect of scattering and converge on a guide star have been shown effective in overcoming scattering [1–16], thus opening an avenue towards high-resolution optical imaging and manipulation in biological tissue. A general concern, however, is that the guide stars may be very weak, leading to the question of how many photons need to be detected to effectively phase conjugate a light wave. Here we demonstrate that phase conjugation is effective even when less than a photon per optical degree of freedom is detected. We found that the best achievable intensity contrast is equal to the total number of detected photons, as long as the resolution of the system is high enough.

Modern optical phase conjugation systems are commonly based on digital optical phase conjugation (DOPC) [10,11,17,18] or use “analog” photorefractive materials [19–21,4]. Both approaches are two-step processes. In the “recording” step, light propagates from a guide star to the phase conjugation system. In DOPC, the phase and amplitude of the scattered wave ( $E^+$ ) is typically measured using phase stepping or off-axis holography using a reference beam. Likewise, in the analog approach, light interferes with a reference beam in order to generate a hologram in the photorefractive material. In the “playback” step, the system generates the phase conjugate copy of the complex field  $E^- \propto (E^+)^*$ , either by using a digital spatial light modulator (SLM) or by illuminating the photorefractive material with a readout beam. Because of the time-reversal symmetry of light propagation, the phase-conjugated wave propagates back through the scattering medium and focuses at the guide star.

In the standard view, phase conjugation increases the intensity in the desired focus by an average factor of  $\eta = M$ , where  $M$  is the number of independently controlled optical modes [1,3,17]. This result assumes that the fields are measured and reproduced exactly. However, until now it

remained unclear how phase conjugation performs in the low-photon limit, where shot noise prevents an accurate measurement of  $E^+$  [22,23]. One may think that the number of detected photons for each measured optical mode needs to be high enough to measure the field in each of the  $M$  modes with sufficient accuracy [24]. However, here we demonstrate the opposite: phase conjugation is still possible even with far less than a single detected photon per degree of freedom. Our study reveals that the fundamental limit for the best possible enhancement  $\eta$  is given by

$$\bar{\eta}_{\max} = \frac{1}{M^{-1} + \bar{n}_r^{-1} + \bar{n}_s^{-1}}, \quad (1)$$

where  $\bar{n}_r$  and  $\bar{n}_s$  are the total number of detected photons coming from the reference and scattered beam, respectively, and the overline denotes averaging over repetitions of the measurement. Interestingly, in the low-photon limit (i.e., when  $\bar{n}_s \ll M$ ), the enhancement is equal to the total number of detected signal photons (i.e.,  $\bar{\eta}_{\max} = \bar{n}_s$ ).

In order to derive Eq. (1), we first introduce the fidelity  $|\gamma|^2$  [25], where

$$\gamma = \frac{\sum_{m=1}^M E_m^- E_m^+}{\sqrt{\sum_{m=1}^M |E_m^-|^2} \sqrt{\sum_{m=1}^M |E_m^+|^2}} \quad (2)$$

with  $m$  the index of the controlled (or measured) optical mode. The fidelity denotes the fraction of the incident power that is shaped correctly and, hence, contributes to the focus. The enhancement and fidelity are related through  $\eta = |\gamma|^2 M$  [25]. Ideally,  $E^- \propto (E^+)^*$ , and  $|\gamma|^2 = 1$ .

In practice, however, shot noise limits the accuracy at which  $E^+$  can be measured. We model the number of detected photons in a single measurement as  $n = \bar{n} + \xi$ , where  $\bar{n}$  is the average photon count and  $\xi$  is a noise term with the statistical properties  $\bar{\xi} = 0$  and  $\overline{\xi^2} = \bar{n}$  following from Poisson statistics [26].

In a phase-stepping interferometry setup, the number of detected photons in a single controlled optical mode is the result of interference between the scattered field  $E_{s,m}$  and the reference field  $E_{r,m}$ :

$$\begin{aligned} n_{k,m} &= |E_{r,m} + E_{s,m}|^2 \\ &= \bar{n}_{k,m} + \xi_{k,m} \\ &= \frac{\bar{n}_{r,m}}{K} + \frac{\bar{n}_{s,m}}{K} + \frac{2}{K} \sqrt{\bar{n}_{r,m}\bar{n}_{s,m}} \cos(\phi_{s,m} + 2\pi k/K) + \xi_{k,m}. \end{aligned} \quad (3)$$

$\xi_{k,m}$  is the shot noise term associated with  $n_{k,m}$  and  $\phi_{s,m}$  is the relative phase difference between the reference field and the scattered field at a chosen mode  $m$ . In our experiments, the controlled mode index  $m$  corresponds to different correlation areas (loosely speaking, different granules in the scattered speckle field; see below for the definition) on the SLM (or image sensor).  $K$  is the total number of phase steps, and  $k \in 1 \dots K$  denotes the step number. For off-axis holography, the relation is identical: to measure the amplitude and phase of a single optical mode,  $K$  different pixels are needed; when we identify these pixels by  $k \in 1 \dots K$ , the same equation is found. The reference intensity is considered to be homogeneous across all optical modes:  $\bar{n}_{r,m} = \bar{n}_r/M$ . The complex amplitude of the signal beam is then reconstructed using

$$E_m^- = \sum_{k=1}^K n_{k,m} e^{-i2\pi k/K}. \quad (4)$$

By substituting Eq. (3) and averaging over measurements, we find

$$\bar{E}_m^- = \sqrt{\bar{n}_{r,m}\bar{n}_{s,m}} e^{i\phi_{s,m}} = \sqrt{\bar{n}_{r,m}} E_m^{+*} \quad (5)$$

and

$$|\bar{E}_m^-|^2 = \bar{n}_{r,m}\bar{n}_{s,m} + \sum_{k=1}^K \bar{\xi}_{k,m}^2, \quad (6)$$

as the dc term  $\bar{n}_{r,m}/K + \bar{n}_{s,m}/K$  and the noise term  $\xi_{k,m}$  in Eq. (3) are eliminated with a summation and averaging, respectively. Then, using  $\bar{\xi}_{k,m}^2 = \bar{n}_{k,m}$  and Eq. (3),

$$|\bar{E}_m^-|^2 = \bar{n}_{r,m}\bar{n}_{s,m} + \bar{n}_{r,m} + \bar{n}_{s,m}. \quad (7)$$

Assuming that  $M$  is sufficiently high, we can write  $\sum_{m=1}^M E_m^- E_m^{+*} \approx \sum_{m=1}^M \bar{E}_m^- E_m^{+*}$  and  $\sum_{m=1}^M |E_m^-|^2 \approx \sum_{m=1}^M |\bar{E}_m^-|^2$ . We can now substitute Eqs. (5) and (7) into Eq. (2) to find

$$\begin{aligned} |\gamma|^2 &= \frac{(\sum_{m=1}^M \bar{n}_{s,m})^2 \bar{n}_r/M}{(\sum_{m=1}^M \bar{n}_{r,m}\bar{n}_{s,m} + \bar{n}_{r,m} + \bar{n}_{s,m})(\sum_{m=1}^M \bar{n}_{s,m})} \\ &= \frac{\bar{n}_s^2 \bar{n}_r/M}{(\bar{n}_s \bar{n}_r/M + \bar{n}_r + \bar{n}_s) \bar{n}_s} = \frac{M^{-1}}{M^{-1} + \bar{n}_r^{-1} + \bar{n}_s^{-1}}, \end{aligned} \quad (8)$$

which gives us the result in Eq. (1). Interestingly, in the low-photon limit, the fidelity equals  $\bar{n}_s/M$ , so increasing the

number of degrees of freedom reduces the fidelity, while the enhancement  $\eta = |\gamma|^2 M$  in the focus remains constant.

*Experiment.*—In our experiment, DOPC [10,11,17,18] is used to generate the phase conjugate beam (see Fig. 1 for a detailed schematic of the experiment). The digital implementation here is particularly useful to quantify the scattered input photon number  $\bar{n}_s$  and to see how it affects the phase conjugation results.

Our experiment is designed to observe the contrast  $\eta$  over a wide range of  $\bar{n}_s$  (see Fig. 1). A 532-nm laser beam is split into three beam paths: input beam, reference beam, and playback beam. All beams are spatially filtered using single mode fibers. Beam splitter BS2 divides the DOPC system into a detection part (CMOS and BS1) and a playback part (SLM and BS3). The SLM is placed in the exact mirror

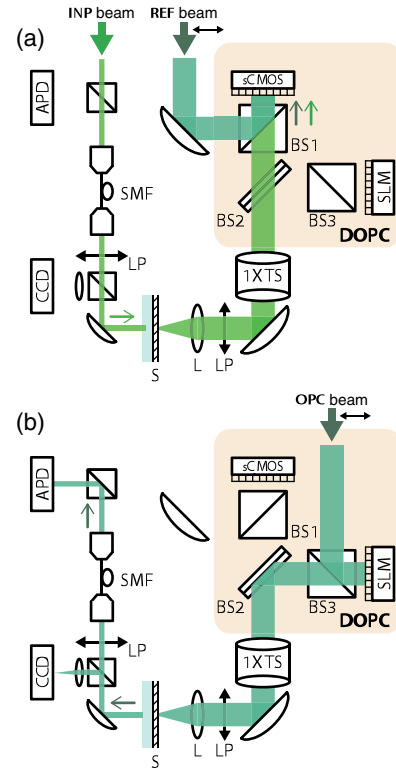


FIG. 1. Experimental setup. (a) In the recording step, the input beam (denoted as “INP beam”), which is spatially filtered through a single-mode fiber (SMF), propagates through the scattering media (S). The scattered input beam is then collected by the planoconvex lens (L, focal length = 5 cm) and relayed to the sCMOS sensor through a 1X telescope system (1X TS), interfering with the collimated reference beam (denoted as “REF beam”). The complex field distribution of the signal beam is then measured using a four-step phase-stepping method. (b) In the playback step, the collimated playback beam (denoted as “OPC beam”) reflects off the spatial light modulator (SLM) on which the conjugated phase of the measured wave front is displayed. The APD measures the fraction of power in the phase conjugate beam that is coupled back to the SMF, and the CCD camera directly captures the transmitted intensity distribution. BS, 50/50 beam splitter; LP, linear polarizer.

image of the image sensor. The single mode fiber on the input beam path defines the input mode, while each correlation area on the SLM (or image sensor) plane defines a controlled (or measured) mode.

In the recording step, the collimated input beam, scattered by a 450- $\mu\text{m}$ -thick opaque opal glass (No. 43-717, Edmund Optics), was interfered with a plane reference beam on the image sensor [PCO.edge 5.5, PCO; 16-bit sCMOS sensor; readout noise  $1.5e^-$  (rms), dark noise  $0.5e^-/s$  per pixel; quantum efficiency  $\sim 50\%$ ]. The number of controlled modes  $M$ , which was experimentally quantified as the total number of pixels ( $1920 \times 1080$ ) divided by the speckle autocorrelation area ( $\sim 9$  pixels) [27], was around 220 000. The speckle correlation area was controlled with an aperture placed in the 1X telescope system which sets the range of transverse wave vector components contributable in the phase conjugation process. Reference photons per controlled optical mode  $\bar{n}_r/M$  were  $\sim 10\,000$ . The average number of signal photons per controlled optical mode  $\bar{n}_s/M$ , which was controlled with a precalibrated acousto-optic modulator (not shown in Fig. 1), was within the range of 0.004–2000. The input wave front is measured using phase-stepping holography with the four phase steps  $0, \pi/2, \pi,$  and  $3\pi/2$ .

In the playback step, the phase conjugate copy of the measured wave front was generated using a phase-only spatial light modulator (PLUTO phase-only SLM, Holoeye). Then, the fraction of power in the phase conjugate beam, coupled back to the single mode fiber through the opal glass, was monitored by the avalanche photodetector (APD). The use of a single mode fiber assures that the optical mode we monitor is in the exact time-reversal symmetry with the input mode. The intensity distribution of the phase conjugate beam was also directly visualized by the CCD sensor.

Even if the incident field could be measured exactly, the accuracy of the reproduction is limited by experimental factors, such as the fact that a phase-only SLM is used instead of the idealized phase-and-amplitude light modulator. Other factors, such as pixel cross talk on the SLM or imperfect alignment further reduce the fidelity. The effect of these experimental factors and the effect of shot noise will be statistically independent. Therefore, we can define an experimental fidelity as  $|\gamma_{\text{ex}}|^2 \equiv \alpha_{\text{ex}}|\gamma|^2$ , where  $|\gamma|^2$  is given by Eq. (8), and the constant  $\alpha_{\text{ex}}$  incorporates all experimental limitations (as discussed in Ref. [28]). For perfect phase-only modulation  $\alpha_{\text{ex}} = \pi/4$  [3], and typical experimental values range from 0.1 to 0.5 [17,18]. The experimental constant  $\alpha_{\text{ex}}$  in our experiment was determined to be 0.27, using a high-power input beam so that the effects of shot noise are negligible.

The experiment confirms our theoretical predictions. In Fig. 2, we plotted the experimental enhancement versus  $\bar{n}_s$  together with the theoretical curve  $\eta \equiv \alpha_{\text{ex}}|\gamma|^2 M$ . The experimental enhancement scales with the total number of

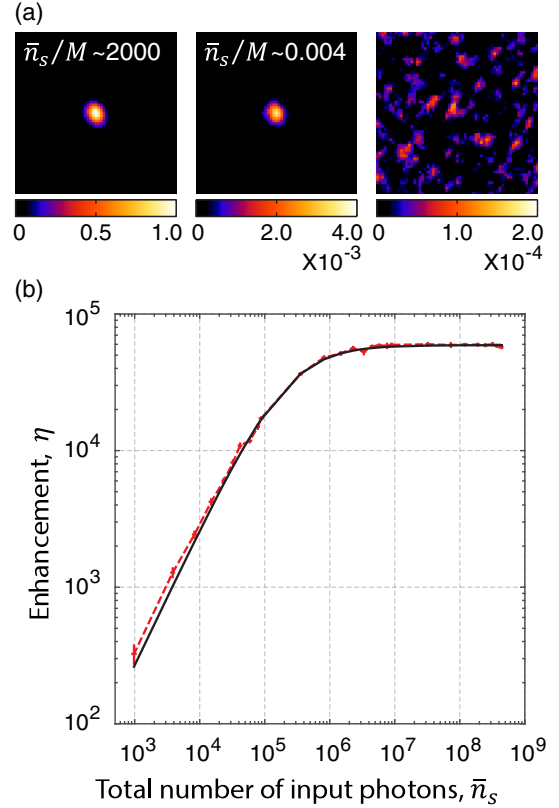


FIG. 2. Effect of shot noise on time-reversal fidelity. (a) The intensity distribution of the phase conjugate beam was directly visualized by the CCD sensor. Left and middle: With the phase conjugate beams generated at  $\bar{n}_s/M$  of  $\sim 2000$  and  $\sim 0.004$ . Right: With an unshaped incident beam (i.e., plane pattern on SLM). (b) Experimental and theoretical enhancement versus the total number of input photons  $\bar{n}_s$ . The solid curve and dotted curve, respectively, represent the theoretical and experimental enhancement.

signal photons  $\bar{n}_s$  exactly as expected over a scaling range of over 5 orders of magnitude. The image captured from the CCD sensor further confirms that the time-reversed mode is reconstructed with a significant contrast corresponding to the scaling relation even with the extremely low value of 0.004 photon per controlled optical mode, corresponding to about 1000 signal photons in total.

In order to have a better understanding of the robustness in the low-photon limit, we further analyzed the influence of shot noise on the phase measurement accuracy for the case that  $K=4$ . By letting  $E_m^- = A_m \cos \phi_m + i A_m \sin \phi_m = [(\bar{n}_{4,m} + \xi_{4,m}) - (\bar{n}_{2,m} + \xi_{2,m})] + i[(\bar{n}_{1,m} + \xi_{1,m}) - (\bar{n}_{3,m} + \xi_{3,m})]$ , the joint probability density function (PDF) for the measured amplitude and phase is given by

$$\begin{aligned}
 P_{A,\phi}(A_m, \phi_m | \bar{n}_{1,m}, \bar{n}_{2,m}, \bar{n}_{3,m}, \bar{n}_{4,m}) \\
 = A_m P_{\Delta}(A_m \cos \phi_m | \bar{n}_{4,m}, \bar{n}_{2,m}) P_{\Delta}(A_m \sin \phi_m | \bar{n}_{1,m}, \bar{n}_{3,m}),
 \end{aligned} \tag{9}$$

as the two differential terms are statistically independent. The leading term  $A_m$  on the right side is the Jacobian of the transformation, and  $P_\Delta(\Delta|\bar{n}_1, \bar{n}_2)$  is the probability mass function of the subtraction of two independent Poissonian random variables with means  $\bar{n}_1$  and  $\bar{n}_2$ , respectively, which is known as the Skellam distribution [29]:

$$P_\Delta(\Delta|\bar{n}_1, \bar{n}_2) = e^{-(\bar{n}_1 + \bar{n}_2)} \left(\frac{\bar{n}_1}{\bar{n}_2}\right)^{\Delta/2} I_\Delta(2\sqrt{\bar{n}_1 \bar{n}_2}), \quad (10)$$

where  $I_\Delta$  is the modified Bessel function of the first kind. To calculate the distribution of the measured phase, we choose an input beam with  $\phi_{s,m} = 0$  and use Eq. (3) to find  $\bar{n}_{k,m}$ . Considering that  $\bar{n}_{s,m}$  of a fully developed speckle pattern is exponentially distributed over different controlled optical modes [27], the PDF for the phase measurement accuracy can be calculated by integrating the PDF given in Eq. (9) over  $\bar{n}_s$  weighted with an exponential PDF. Finally, we integrate over all possible measured amplitudes  $A$ :

$$\begin{aligned} P_\phi(\phi) = & \iint \frac{e^{-\bar{n}_{s,m}/\langle \bar{n}_{s,m} \rangle}}{\langle \bar{n}_{s,m} \rangle} A_m P_\Delta \left( A_m \cos \phi_m \left| \frac{\bar{n}_{r,m}}{4} + \frac{\bar{n}_{s,m}}{4} \right. \right. \\ & \left. \left. + \frac{\sqrt{\bar{n}_{r,m} \bar{n}_{s,m}}}{2}, \frac{\bar{n}_{r,m}}{4} + \frac{\bar{n}_{s,m}}{4} - \frac{\sqrt{\bar{n}_{r,m} \bar{n}_{s,m}}}{2} \right) \right. \\ & \times P_\Delta \left( A_m \sin \phi_m \left| \frac{\bar{n}_{r,m}}{4} + \frac{\bar{n}_{s,m}}{4}, \frac{\bar{n}_{r,m}}{4} + \frac{\bar{n}_{s,m}}{4} \right) \right) \\ & \times dA_m d\bar{n}_{s,m}, \end{aligned} \quad (11)$$

where  $\langle \bar{n}_{s,m} \rangle$  ( $= \bar{n}_s/M$ ) denotes the average signal strength in a single controlled optical mode.

Figure 3 presents the excellent agreement between the experimental phase error distribution  $P_\phi(\Delta\phi)$  and the theoretical distribution [Eq. (11), evaluated numerically]. The experimental phase error distribution is obtained by comparing the phase map measured at the low-photon limit with the reference phase map measured at  $\bar{n}_s/M \sim 2000$ , where the impact of shot noise is negligible. The phase errors are binned into 101 equally divided bins. As expected, the phase measurement accuracy is significantly reduced as the number of signal photons decreases below the level of  $\bar{n}_s/M \sim 1$ . However, the phase error is favorably populated around 0 even down to  $\bar{n}_s/M \sim 0.004$ . The excellent statistical stability of the experimental results can be attributed to the large number of measured modes ( $\sim 220\,000$ ), which itself yields an ensemble-averaged result.

The counterintuitive aspects of our experiments can best be illustrated by the following thought experiment. Imagine a stream of 1000 signal photons detected and played back by  $10^7$  controllable modes. In the quantized picture that at least a single photon is required for a correct phase measurement, we would at best make a correct measurement at  $M = 10^3$  modes, which is  $1/10^4$  of the controlled modes (i.e.,  $|\gamma| = 1/10^4$ ), resulting in the

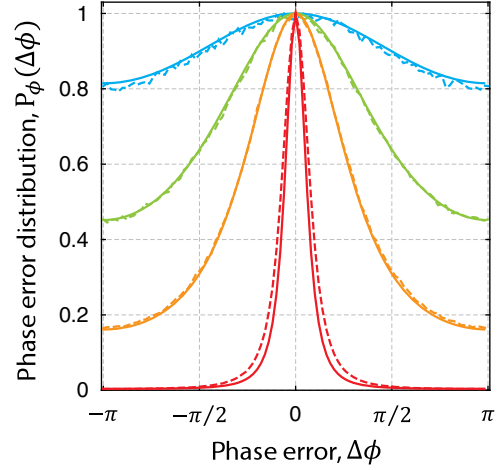


FIG. 3. Effect of shot noise on phase measurement accuracy. The phase error distribution  $P_\phi(\Delta\phi)$  was obtained by comparing the phase map measured at a low-photon limit (red curve,  $\bar{n}_s/M \sim 20$ ; orange curve,  $\bar{n}_s/M \sim 0.4$ ; green curve,  $\bar{n}_s/M \sim 0.07$ ; cyan curve,  $\bar{n}_s/M \sim 0.004$ ) with the reference phase map measured at a negligible shot noise level ( $\bar{n}_s/M \sim 2000$ ). The number of bins is 101, uniformly distributed from  $-\pi$  to  $\pi$ . The solid curve and dotted curve, respectively, represent the theoretical and experimental phase error distribution. All curves are normalized such that the maximum value is unity.

unappreciable enhancement of  $1 + 10^7/10^8 = 1.1$  [3]. This prediction is at odds with our presented analysis and experimental results that promise an enhancement of 1000. This paradox can be explained by the concept of wave-particle duality—the quantization occurs after the interferometric mixing. Therefore, even though each DOPC pixel harvests wave front information from less than a quantum of signal photon energy on average, the information is still useful in steering the output phasor, such that in aggregate we can achieve a significant contrast in time-reversed light focusing through a scattering medium.

For a given number of signal photons, the contrast that can be achieved with optical phase conjugation is much higher than what can be achieved with feedback-based methods. For feedback-based wave front shaping,  $M$  different interferometric measurements are needed, and the maximum enhancement is proportional to the average number of photons recorded per measurement, i.e.,  $\eta_{\max} = \bar{n}_s/M$ , with  $\bar{n}_s$  the total photon budget for all measurements [30]. Note, however, that feedback-based wave front shaping allows one to simultaneously determine the wave fronts for focusing at any desired point [31], whereas optical phase conjugation finds only the solution for a single focus at the time. Therefore, for measuring the full transmission matrix of a sample (with all possible combinations of  $M$  input modes and  $M$  output modes), both methods are expected to reach the same fidelity for a given total photon budget.

In summary, our study provides the first comprehensive understanding of the interplay between the number of controllable modes and photon budgets available in digital

optical phase conjugation. As indicated in Eq. (1), in general, having one of the three parameters ( $M$ ,  $\bar{n}_r$ , and  $\bar{n}_s$ ) much larger than the other does not significantly improve the enhancement  $\eta$ . In other words, the enhancement  $\eta$  is capped to the least of three parameters.

Our results, which are directly applicable to the conventional analog phase conjugation scheme [19–21,4], exhibit numerous counterintuitive implications:

(i) It is always beneficial to perform phase conjugation with a resolution as high as possible, even if doing so decreases the number of photons per controlled mode to far below 1. Our results guarantee that this configuration provides an enhancement equal to the total number of photons  $\bar{n}_s$  collected during the phase conjugation process.

(ii) In the low-photon limit, the enhancement, traditionally known to solely depend on the number of controlled optical modes, becomes proportional to the parameters affecting  $\bar{n}_s$ , such as the incident power of the input beam, the exposure time for wave front measurement, and the tagging efficiency of the guide star (e.g., diffraction efficiency of the ultrasound guide star [6–8] and quantum yield of fluorescence guide star [9]).

(iii) The number of input modes [i.e., the number of optical modes in guide star(s)] can be increased without compromising the enhancement. In the low-photon limit, the effect of the increased photon number  $\bar{n}_s$ , that results from increasing the number of input modes, will counteract the well-known inverse scaling of the enhancement with the number of input modes [3].

Our results demonstrate that phase conjugation can be used even when the photon budget is extremely low, such as in high-speed focusing through dynamic media or imaging deep inside tissue. We anticipate the analysis and implications, presented in this study, will serve as a stepping stone towards the adoption and improvement of optical phase conjugation in realizing its biomedical applications.

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